

Infinite dimensional Ornstein-Uhlenbeck processes driven by Lévy processes

David Applebaum

*School of Mathematics and Statistics
University of Sheffield
Hicks Building, Hounsfield Road
Sheffield, England, S3 7RH
e-mail: D.Applebaum@sheffield.ac.uk*

Abstract: We review the probabilistic properties of Ornstein-Uhlenbeck processes in Hilbert spaces driven by Lévy processes. The emphasis is on the different contexts in which these processes arise, such as stochastic partial differential equations, continuous-state branching processes, generalised Mehler semigroups and operator self-decomposable distributions. We also examine generalisations to the case where the driving noise is cylindrical.

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1. Introduction

The physical phenomenon of Brownian motion is well described by the stochastic process that bears its name, which originated in the work of Einstein, Bachelier, Smoluchowski and Wiener, provided the particle moves “freely” in its surround-

ing medium, subject only to the random force of molecular bombardment. If the viscous properties of the medium are also taken into account, Ornstein and Uhlenbeck [37] proposed that the velocity $v(t)$ at time t of a diffusing particle of mass m should satisfy the Langevin equation:

$$m \frac{dv}{dt} + \beta v = F(t), \quad (1.1)$$

where β is the viscosity coefficient, and F is a random force acting on the particle. Doob [23] rewrote (1.1) as an SDE in which the formal differential $F(t)dt$ was replaced by the stochastic differential $dB(t)$ of Brownian motion. Its solution (with $m = 1$, for convenience) is the prototypical *Ornstein-Uhlenbeck process*,

$$v(t) = e^{-\beta t} v(0) + \int_0^t e^{-\beta(t-s)} dB(s),$$

wherein the right-hand-side is the standard Wiener stochastic integral, soon to be extensively generalised by Itô. The process $(v(t), t \geq 0)$ is both Gaussian and Markov, and if we choose $v(0)$ to have the limiting distribution $N(0, 1/2\beta)$, then it is also stationary.

From a probabilistic point of view, it is natural to consider different driving noises, and also to work in a multivariate framework. A number of authors have investigated the case where B is generalised to be an arbitrary Lévy process, and the corresponding generalised Ornstein-Uhlenbeck processes are interesting from both a theoretical and applied viewpoint (see e.g. [5] pp. 241–243 and references therein). Most of the theoretically interesting properties in the finite-dimensional case, such as the structural relationship with the class of self-decomposable distributions, will re-emerge in the infinite-dimensional framework that we consider below; among the most important areas of application we mention volatility modelling in mathematical finance [10, 31], and physical models of anomalous diffusion [27].

In this review, we will be concerned solely with infinite-dimensional generalisations where the noise is a general Lévy process taking values in a real Hilbert space. The rationale for this choice, is that this class of processes is very rich, and that they lie at the intersection of a number of interesting areas in probability theory and stochastic analysis. In particular this article will consider their role in the study of stochastic partial differential equations (SPDEs), generalised Mehler semigroups, operator self-decomposable distributions, continuous-state branching processes, and cylindrical measures and processes. Some of these connections have been investigated by other authors, including the current one, see e.g. [3, 11, 34, 29], but it seems to be a useful endeavour to discuss them all within a single article, and to give readers some pointers to the large and growing literature in these active areas of research.

We will not consider the case herein where the noise is purely Gaussian. This is because that theory has been very well developed, and the interested reader can find full accounts in the monographs [20, 22], and references therein. We also resisted the temptation to generalise further to Banach space-valued processes, as that theory has not yet reached such a mature stage as the Hilbert space one.

The plan of the paper is as follows. After some preliminaries on Lévy processes and stochastic integrals in section 2, we define infinite-dimensional Ornstein-Uhlenbeck processes in section 3, and see how they arise naturally in the study of SPDEs with additive Lévy noise. In section 4, we collect together some of the basic probabilistic properties of these processes. Mehler semigroups of operators, skew-convolution semigroups of measures and their generalisations are the topic of section 5. In section 6, we turn to invariant measures and operator self-decomposability. Finally in section 7, we look at various approaches to defining Ornstein-Uhlenbeck processes driven by cylindrical Lévy noise.

Preliminaries and notation. Throughout this work H is a real separable Hilbert space, with inner product $\langle \cdot, \cdot \rangle$ and norm $\|\cdot\|$, $\mathcal{B}(H)$ is its Borel σ -algebra, and $\mathcal{M}_1(H)$ is the space of all Borel probability measures defined on H . We denote by $B_b(H)$ the Banach space (with respect to the supremum norm) of all real-valued, bounded Borel measurable functions defined on H , and by $C_b(H)$, its closed subspace of bounded continuous functions. The open ball of radius 1 in H , centered at the origin, will always be denoted by B_1 . A positive self-adjoint linear operator T in H is *trace-class* if its trace $\text{tr}(T) = \sum_{n=1}^{\infty} \langle T e_n, e_n \rangle < \infty$, for some (and hence all) complete orthonormal basis $(e_n, n \in \mathbb{N})$ in H . The algebra of all bounded linear operators on H will be denoted by $L(H)$.

The *Fourier transform* (or characteristic function) of $\mu \in \mathcal{M}_1(H)$ is the bounded continuous mapping $\widehat{\mu} : H \rightarrow \mathbb{C}$ defined for all $u \in H$ by

$$\widehat{\mu}(u) = \int_H e^{i\langle u, x \rangle} \mu(dx).$$

The *convolution* $\mu_1 * \mu_2$ of $\mu_1, \mu_2 \in \mathcal{M}_1(H)$, is the unique Borel probability measure for which

$$\int_H f(x) (\mu_1 * \mu_2)(dx) = \int_H \int_H f(x+y) \mu_1(dx) \mu_2(dy),$$

for all $f \in B_b(H)$. Note that $\mu_1 * \mu_2 = \mu_2 * \mu_1$, and for all $u \in H$,

$$\widehat{\mu_1 * \mu_2}(u) = \widehat{\mu_1}(u) \widehat{\mu_2}(u).$$

Suppose that Q is a probability kernel, i.e. a mapping $H \times \mathcal{B}(H) \rightarrow [0, 1]$, so that for fixed $A \in \mathcal{B}(H)$, the mapping $x \rightarrow Q(x, A)$ is measurable, and for fixed $x \in H$, $Q(x, \cdot) \in \mathcal{M}_1(H)$. If $\rho \in \mathcal{M}_1(H)$, then $Q\rho$ is defined to be the unique measure in $\mathcal{M}_1(H)$ for which

$$\int_H f(x) (Q\rho)(dx) = \int_H \int_H f(y) Q(x, dy) \rho(dx),$$

for all $f \in B_b(H)$. If $T : H \rightarrow H$ is a bounded linear operator, and $Q(x, \cdot) = \delta_{Tx}$ (where δ_y is the usual Dirac mass at $y \in H$), then $Q\rho = \rho \circ T^{-1}$, which we also write succinctly as $T\rho$.

2. Lévy processes and stochastic integration in Hilbert space

Let (Ω, \mathcal{F}, P) be a probability space that is equipped with a filtration $(\mathcal{F}_t, t \geq 0)$, which satisfies the “usual hypotheses” of right-continuity and completeness. Let H be a real separable Hilbert space. A *Lévy process* is an adapted H -valued process $L = (L(t), t \geq 0)$ for which

- $L(0) = 0$ (a.s.),
- L is stochastically continuous,
- L has stationary and independent increments, where the latter is in the strong sense that $L(t) - L(s)$ is independent of \mathcal{F}_s , for all $0 \leq s < t < \infty$,
- L has càdlàg paths.

Then L has the following Lévy-Itô decomposition (see [1, 48]), for all $t \geq 0$:

$$L(t) = bt + B_Q(t) + \int_{\|x\| \leq 1} x \tilde{N}(t, dx) + \int_{\|x\| > 1} x N(t, dx), \quad (2.2)$$

wherein

- the vector $b \in H$,
- the process $B_Q = (B_Q(t), t \geq 0)$ is a Brownian motion taking values in H with covariance operator Q , so that for all $s, t \geq 0, \psi, \phi \in H$,

$$\mathbb{E}(\langle \psi, B_Q(s) \rangle \langle \phi, B_Q(t) \rangle) = \langle Q\psi, \phi \rangle (s \wedge t),$$

where Q is a positive, symmetric, trace-class operator,

- N is a Poisson random measure on $[0, \infty) \times H$, which is independent of B_Q , and has compensator:

$$\tilde{N}(dt, dx) = N(dt, dx) - dt\nu(dx),$$

where ν is a Lévy measure on H , i.e.

$$\nu(\{0\}) = 0 \text{ and } \int_H (\|x\|^2 \wedge 1) \nu(dx) < \infty.$$

From this we can deduce the Lévy-Khintchine formula, for all $t \geq 0, u \in H$,

$$\mathbb{E}(e^{i\langle u, L(t) \rangle}) = e^{t\eta(u)},$$

where the *characteristic exponent* $\eta : H \rightarrow \mathbb{C}$ is given by

$$\eta(u) = i\langle u, b \rangle - \frac{1}{2} \langle Qu, u \rangle + \int_H (e^{i\langle u, y \rangle} - 1 - i\langle u, y \rangle \mathbf{1}_{B_1}(y)) \nu(dy). \quad (2.3)$$

The law of $L(t)$ is uniquely determined by its *characteristics* (b, Q, ν) .

The process L is a martingale, and will be called a *Lévy martingale*, if and only if $\int_{|x|>1} |x| \nu(dx) < \infty$ and $b = -\int_{|x|>1} x \nu(dx)$ (see e.g. [5] p. 133) in which case we may write (2.2) as

$$L(t) = B_Q(t) + \int_H x \tilde{N}(t, dx).$$

If $L(t) = L_M(t) + bt$, where L_M is a Lévy martingale, we call L a *Lévy martingale with drift*.

If $F = (F(t), t \geq 0)$ is a suitably regular adapted, bounded operator-valued process in H , we can make sense of stochastic integrals $\int_0^t F(t)dL(t)$, either by using the general approach of [36] (see in particular, sections 1.2 and 6.14), or by using the Lévy-Itô decomposition (2.2), as in [3] and Chapter 8 of [42], and defining:

$$\begin{aligned} \int_0^t F(s)dL(s) &= \int_0^t F(s)b \, ds + \int_0^t F(s)dB_Q(s) \\ &+ \int_0^t \int_{\|x\| \leq 1} F(s)x\tilde{N}(s, dx) + \int_0^t \int_{\|x\| > 1} F(s)xN(s, dx). \end{aligned}$$

Here the first integral is a Bochner integral, and the last is a random sum. The middle two may be defined by using stochastic integration against martingale-valued measures. In this paper we will only be concerned with the case where F is deterministic. In that case, we may also define $\int_0^t F(s)dL(s)$ using integration by parts, as in [30].

Note. A Borel probability measure on H is infinitely divisible if and only if its characteristic exponent is given by (2.3); see Theorem 4.10, pp. 181–182 in [40].

3. Ornstein-Uhlenbeck processes and stochastic partial differential equations

Let $(S(t), t \geq 0)$ be a C_0 -semigroup acting in H (i.e. a strongly continuous one-parameter semigroup of bounded, linear operators on H). Let Y_0 be a fixed \mathcal{F}_0 -measurable random variable, and L be an H -valued Lévy process. The *Ornstein-Uhlenbeck process* (or OU process, for short) associated with these data is the adapted process $Y = (Y(t), t \geq 0)$ defined for $t \geq 0$ by

$$Y(t) = S(t)Y_0 + \int_0^t S(t-r)dL(r). \quad (3.4)$$

Observe that if $H = \mathbb{R}$, then we must have $S(t) = e^{\lambda t}$ for some $\lambda \in \mathbb{R}$, and we recapture the usual Lévy-driven one-dimensional Ornstein-Uhlenbeck process (see e.g. [5] pp. 237–243, and references therein). When H is finite dimensional we may consider matrix semigroups of the form $S(t) = e^{tA} = \sum_{n=0}^{\infty} \frac{t^n A^n}{n!}$. Such objects are discussed from a distributional perspective in [50] pp. 106–114. From now on we will always assume that H is infinite-dimensional.

The stochastic integral in (3.4) always make sense using the procedures that we have sketched (see also [15]). The process $t \rightarrow \int_0^t S(t-r)dL(r)$ is called a *stochastic convolution*. It is not, in general, a semimartingale; and in contrast to the finite-dimensional case, when L is a Lévy martingale, the stochastic convolution is not the product of a martingale with a deterministic process. This is because we cannot decompose the semigroup as “ $S(t-r) = S(t)S(-r)$ ”. This causes problems in studying the time-regularity of the OU process, which will be discussed in section 4.

Motivation for studying the OU process comes from the study of *stochastic partial differential equations*, or SPDEs for short, that are driven by Lévy noise. Consider, for example the heat equation in \mathbb{R}^d ,

$$\frac{\partial u}{\partial t} = \sum_{j=1}^d \frac{\partial^2 u}{\partial x_j^2}, \quad (3.5)$$

with initial condition $u(0) = u_0 \in L^2(\mathbb{R}^d)$. One approach to randomizing this equation would be to introduce a suitable two-parameter Lévy sheet $(L(x, t), x \in \mathbb{R}^d, t \geq 0)$ and seek to give meaning to the equation

$$\frac{\partial Y(t, x)}{\partial t} = \sum_{j=1}^d \frac{\partial^2 Y(t, x)}{\partial x_j^2} + \dot{L}(x, t),$$

whose solution would be a random field on $[0, \infty) \times \mathbb{R}^d$.

This procedure is implemented in [9], by extending the methodology developed by Walsh [54] for the case of Gaussian noise. An alternative approach is to recognise that the Laplacian $\Delta = \sum_{j=1}^d \frac{\partial^2}{\partial x_j^2}$ is the infinitesimal generator of the heat semigroup $(S_h(t), t \geq 0)$ in $L^2(\mathbb{R}^d)$. Then the solution of the heat equation (3.5) is given by $u(t) = S_h(t)u_0$ in $L^2(\mathbb{R}^d)$. From this point of view, we may introduce Lévy noise by using the Hilbert space-valued Lévy process L and write the SPDE as an infinite-dimensional, Itô-sense stochastic differential equation

$$dY(t) = \Delta Y(t)dt + dL(t). \quad (3.6)$$

Next we generalise, and appreciate that there is nothing special from the point of view of Itô calculus, about the semigroup $(S_h(t), t \geq 0)$ acting in $L^2(\mathbb{R}^d)$. In fact we may consider an arbitrary C_0 -semigroup $(S(t), t \geq 0)$ having infinitesimal generator A , and acting in a Hilbert space H . Then we seek to make sense of solutions to:

$$dY(t) = AY(t)dt + dL(t). \quad (3.7)$$

We regard equation (3.7) as an infinite-dimensional *Langevin equation*. It has a unique mild solution which is given by the variation of constants formula, and this is precisely the OU process (3.4). As is shown in Theorems 9.15 and 9.29 of [42], or in [3], this is also the unique weak solution in that for all $u \in \text{Dom}(A^*)$, $t \geq 0$, with probability 1,

$$\langle u, Y(t) - Y_0 \rangle = \langle u, L(t) \rangle + \int_0^t \langle A^* u, Y(s) \rangle ds. \quad (3.8)$$

Equation (3.7) is the simplest SPDE with additive Lévy noise. More generally, one may consider more complicated equations with multiplicative noise of the form:

$$dY(t) = (AY(t) + F(t, Y(t-)))dt + G(t, Y(t-))dL(t). \quad (3.9)$$

For a comprehensive introduction to the study of these equations, see [42]. The equation (3.9) is called *semilinear* if $G(t, Y(t-)) = G$ is a bounded linear

operator, and in this case in particular, knowledge of the case $F \equiv 0$ (i.e. the Ornstein-Uhlenbeck process), may be an important first step in obtaining information about the solution to the more general equation, see e.g. [44, 2]. The two different approaches to solving SPDEs that we have presented above, using on the one hand, two-parameter “space-time white noise”, and on the other hand, infinite-dimensional processes, have each generated a considerable literature; nonetheless there are important cases where they both give rise to the same solution, see e.g. [16].

Note. A rather straightforward generalisation of (3.7) considers equations of the form

$$dY(t) = AY(t)dt + BdL(t),$$

where L takes values in a different Hilbert space H_1 , and B is a bounded linear operator from H_1 to H . Since $(BL(t), t \geq 0)$ is itself a Lévy process in H , this is already included in the case that we consider.

4. Basic properties of Ornstein-Uhlenbeck processes

It follows from elementary properties of stochastic integrals with respect to Lévy processes that $(Y(t) - Y_0, t \geq 0)$ is an additive process, i.e. a stochastically continuous process, which vanishes at zero (a.s.), and has independent increments. Hence, for $t > 0$, the random variable $Y(t) - Y_0$ is infinitely divisible. From our assumptions, $Y(t) - Y_0$ is independent of Y_0 , and so, if Y_0 is infinitely divisible, then so is $Y(t)$. If Y_0 is infinitely divisible with characteristic exponent η_0 , we have (see [15, 4]), for all $u \in H$,

$$\mathbb{E}(e^{i\langle u, Y(t) \rangle}) = \exp \left\{ \eta_0(S(t)^*(u)) + \int_0^t \eta(S(r)^*(u)) dr \right\}.$$

Furthermore, if Y_0 has characteristics (b_0, Q_0, ν_0) , then $Y(t)$ has characteristics (b_t, Q_t, ν_t) where:

$$\begin{aligned} b_t &= S(t)b_0 + \int_H S(t)x[\mathbf{1}_{B_1}(S(t)x) - \mathbf{1}_{B_1}(x)]\nu_0(dx) \\ &+ \int_0^t S(r)bdr + \int_0^t \int_H S(r)x[\mathbf{1}_{B_1}(S(r)x) - \mathbf{1}_{B_1}(x)]\nu(dx)dr, \\ Q_t &= S(t)Q_0S(t)^* + \int_0^t S(r)QS(r)^*dr, \\ \nu_t(B) &= \nu_0(S(t)^{-1}(B)) + \int_0^t \nu_0(S(r)^{-1}(B))dr, \end{aligned}$$

for each $B \in \mathcal{B}(H)$.

In the frequently encountered case where $Y_0 = y_0$ (a.s.), with fixed $y_0 \in H$, then $\eta_0(S(t)^*u) = \langle u, S(t)y_0 \rangle$.

Necessary and sufficient conditions for the OU process Y to have càdlàg paths have not yet been established, and this seems to be a difficult problem, in general. Two sufficient conditions are described in [42], pp. 156–161. In both cases we require the process L to be a square-integrable Lévy martingale with drift.

1. *The Kotelenez approach.* This utilises a generalisation of Doob’s submartingale inequality to stochastic convolution with respect to martingales. In our case, the conclusion is that the stochastic convolution (and hence the OU process) has a càdlàg version if $(S(t), t \geq 0)$ is a generalised contraction semigroup, i.e. there exists $\beta > 0$ so that $\|S(t)\| \leq e^{\beta t}$ for all $t \geq 0$.
2. *The Hausenblas-Seidler approach.* This requires $(S(t), t \geq 0)$ to be a contraction semigroup. It utilises the Sz-Nagy–Foias theory of dilations of such semigroups to embed H as a closed subspace of a larger Hilbert space H_1 , and assert that there is a strongly continuous unitary group $(U(t), t \in \mathbb{R})$ in that larger space so that $S(t) = PU(t)$ for all $t \geq 0$, where P is the orthogonal projection of H_1 onto H . Then the stochastic convolution has a càdlàg version. A key feature of the proof is to work in H_1 , and to use the group property to write

$$\int_0^t U(t-s)dL(s) = U(t) \int_0^t U(-s)dL(s).$$

Since it is the solution to an SDE, $(Y(t), t \geq 0)$ is a Markov process, and we may compute the action of the corresponding transition semigroup $(P_t, t \geq 0)$ on the space $B_b(H)$ from (3.4) to obtain for each $t \geq 0, f \in B_b(H), x \in H$,

$$\begin{aligned} P_t f(x) &= \mathbb{E}(f(Y(t)) | Y_0 = x) \\ &= \int_0^t f(S(t)x + y) \mu_t(dy), \end{aligned} \tag{4.10}$$

where μ_t is the law of $\int_0^t S(t-r)dL(r)$, or equivalently, by stationary increments of L , the law of $\int_0^t S(r)dL(r)$. Note that the semigroup $(P_t, t \geq 0)$ is not, in general, strongly continuous on $C_b(H)$ or its subspace of bounded, uniformly continuous real-valued functions on H ; this is discussed for the Gaussian case in [22], p. 111, and the same argument works in the Lévy case. It is however always quasi-equicontinuous with respect to a weaker topology, called the mixed topology, in $C_b(H)$. This is the strongest locally convex topology, which agrees on bounded sets with the topology of uniform convergence on compacta. The semigroup $(P_t, t \geq 0)$ then has an infinitesimal generator \mathcal{L} (in this generalised sense), which has the following representation for functions f in a certain dense domain (see [4] for details):

$$\begin{aligned} \mathcal{L}f(x) &= \langle A^*(Df)(x), x \rangle + \langle (Df)(x), b \rangle + \frac{1}{2} \text{tr}((D^2f)(x)Q) \\ &+ \int_H [f(x+y) - f(x) - \langle (Df)(x), y \rangle \mathbf{1}_{B_1}(y)] \nu(dy), \end{aligned}$$

where $x \in H$, and D is the Fréchet derivative. Pseudo-differential operator representations of \mathcal{L} have been obtained in [32], and in Proposition 4.1 of [4].

The measures $(\mu_t, t \geq 0)$ have an interesting property, which will play an important role in the sequel. Observe that for all $r, t \geq 0$:

$$\int_0^{r+t} S(v)dL(v) = \int_0^r S(v)dL(v) + \int_r^{r+t} S(v)dL(v),$$

and that the two stochastic integrals on the right hand side of this equation are independent. Furthermore, by stationary increments of L , and the semigroup property we have

$$\int_r^{r+t} S(v)dL(v) \stackrel{d}{=} \int_0^t S(r+v)dL(v) = S(r) \int_0^t S(v)dL(v).$$

We conclude that

$$\mu_{r+t} = \mu_r * S(r)\mu_t. \quad (4.11)$$

It is also not difficult to verify that for all $u \in H, t \geq 0$,

$$\widehat{\mu}_t(u) = \exp \left\{ \int_0^t \eta(S(r)^*u) dr \right\}. \quad (4.12)$$

5. Mehler semigroups and skew-convolution semigroups

In this section we look at a more abstract approach to some of the ideas that were encountered in the last section. Let us suppose that we are given, as above, a C_0 -semigroup $(S(t), t \geq 0)$ on H , and also a family $(\rho_t, t \geq 0)$ of Borel probability measures on H , with $\rho_0 = \delta_0$. Define the linear operators $(T_t, t \geq 0)$ on $B_b(H)$ by

$$T_t f(x) = \int_0^t f(S(t)x + y)\rho_t(dy), \quad (5.13)$$

for each $t \geq 0, f \in B_b(H), x \in H$.

A natural question to ask is, when does $(T_t, t \geq 0)$ satisfy the semigroup property, i.e. $T_{s+t} = T_s T_t$, for all $s, t \geq 0$? As is shown in Proposition 2.2 of [11], this will hold if and only if $(\rho_t, t \geq 0)$ is a *skew-convolution semigroup*, i.e. for all $r, t \geq 0$,

$$\rho_{r+t} = \rho_r * S(r)\rho_t, \quad (5.14)$$

or equivalently, for each $u \in H$,

$$\widehat{\rho_{r+t}}(u) = \widehat{\rho}_r(u)\widehat{\rho}_t(S(r)^*u).$$

When (5.14) holds, we call $(T_t, t \geq 0)$ a (generalised) *Mehler semigroup* as (5.13) generalises the classical Mehler formula for one-dimensional Gaussian Ornstein-Uhlenbeck processes (see e.g. [5], p. 405). Mehler semigroups have been extensively studied in recent years, as analytic objects in their own right, particularly by M. Röckner and his collaborators, and we draw the reader's attention to [11, 25, 32, 33, 49, 39]. Skew convolution semigroups have appeared in other contexts (see below) and are also called *measure-valued cocycles* – see e.g. [29].

We have seen in the last section, that Mehler semigroups appear as the transition semigroups of Hilbert space-valued OU processes. We might ask if the converse is valid, and if, starting from a Mehler semigroup, we can construct a càdlàg OU process for which it is the transition semigroup. The answer is affirmative and we will sketch the idea of the proof, drawing on the accounts in [11, 25], and section 11.5 of [34].

Following [25], we assume that for all $u \in H$, the mapping $t \rightarrow \hat{\rho}_t(u)$ is absolutely continuous on $[0, \infty)$, differentiable at $t = 0$, and defining $\eta(u) = \frac{d}{dt} \hat{\rho}_t(u)|_{t=0}$, that $t \rightarrow \eta(S(t)^*u)$ is locally integrable. It follows that (c.f. (4.12))

$$\hat{\rho}_t(u) = \exp \left\{ \int_0^t \eta(S(r)^*u) dr \right\}. \quad (5.15)$$

Furthermore the mapping η is negative definite, and we may define a weakly-continuous convolution semigroup of probability measures $(\xi_t, t \geq 0)$ on H so that for all $t \geq 0$,

$$\hat{\xi}_t(u) = e^{t\eta(u)}.$$

Then we may obtain a càdlàg Lévy process $L = (L(t), t \geq 0)$ defined on some probability space, and taking values in H , having characteristic exponent η by using standard techniques (see e.g. Chapters 1 and 2 of [5]).

The next step is to construct a Hilbert space \tilde{H} , in which H is continuously embedded, so that $(S(t), t \geq 0)$ extends to a C_0 -semigroup $(\tilde{S}(t), t \geq 0)$ in \tilde{H} whose generator \tilde{A} is an extension of A , with $H \subset \text{Dom}(\tilde{A})$. To see how to do this, we first note that there exists $M \geq 1$ and $c \geq 0$ so that for all $t \geq 0$, $\|S(t)\| \leq Me^{tc}$. Choose $\lambda > c$ and let $R_\lambda = (\lambda I - A)^{-1}$ be the corresponding resolvent of A . We define a new inner product $\langle \cdot, \cdot \rangle_1$ on H by the prescription

$$\langle \phi, \psi \rangle_1 = \langle R_\lambda \phi, R_\lambda \psi \rangle,$$

and define \tilde{H} to be the completion of H in the corresponding norm. Now for all $\tilde{u} \in \tilde{H}$, we define

$$Y(t) = \tilde{S}(t)\tilde{u} + L(t) + \int_0^t \tilde{S}(t-s)\tilde{A}L(s)ds. \quad (5.16)$$

By formal differentiation of (5.16), we see that (3.4) is satisfied by this process (with respect to the extended Hilbert space, and semigroup), with initial condition $Y(0) = \tilde{u}$. Observe that the construction of \tilde{H} ensures that the integral in (5.16) makes sense as $L(s) \subseteq \text{Dom}(\tilde{A})$ for all $s \geq 0$. As pointed out by Li in [34], we may take $\tilde{u} \in H$ to get an H -valued version of the process, but this may not have càdlàg paths.

In the case where the skew-convolution semigroup is square-integrable, so that $\int_H \|x\|^2 \rho_t(dx) < \infty$ for all $t \geq 0$, it is shown in [34], section 11.4. that the regularity conditions assumed above to obtain (5.15) may be dropped; but the price we pay for this is that the measures $(\xi_t, t \geq 0)$ will live on a larger Hilbert space than H , consisting of the locally square-integrable entrance paths for the semigroup $(S(t), t \geq 0)$.

We now discuss an interesting extension of the notion of skew convolution semigroup, which is due to Dawson and Li [17]. Let $(E, +)$ be an abelian Haus-

dorff semigroup, and for $t \geq 0$, let $Q_t : E \times \mathcal{B}(E) \rightarrow [0, 1]$ be the transition kernels of a Markov process. We assume that the underlying process has a *branching property* so that for all $t \geq 0, x_1, x_2 \in E$

$$Q_t(x_1 + x_2, \cdot) = Q_t(x_1, \cdot) * Q_t(x_2, \cdot). \quad (5.17)$$

We say that a family of probability measures $(\mu_t, t \geq 0)$ on $(E, \mathcal{B}(E))$ is a *generalised skew-convolution semigroup* if for all $s, t \geq 0$,

$$\mu_{s+t} = Q_t \mu_s * \mu_t. \quad (5.18)$$

Observe that (5.18) coincides with (5.14), when we take $E = H$ and $Q_t(x, \cdot) = \delta_{S(t)x}(\cdot)$. Now define for $t \geq 0, x \in E$, the kernels

$$Q_t^\mu(x, \cdot) = Q_t(x, \cdot) * \mu_t. \quad (5.19)$$

The next result gives an extension of the Mehler semigroup/OU process circle of ideas to this more general context. We include the proof for the reader's convenience, as we haven't found one in the literature.

Theorem 5.1. $(Q_t^\mu, t \geq 0)$ is the transition kernel of a Markov process.

Proof. We must show that the Chapman-Kolomogorov equations are satisfied, i.e. that for all $f \in B_b(E), x \in E$,

$$\int_S f(z) Q_{s+t}^\mu(x, dz) = \int_S \int_S f(z) Q_t^\mu(y, dz) Q_s^\mu(x, dy).$$

We begin with the right hand side of the last display and use (5.19) twice to obtain

$$\begin{aligned} & \int_S \int_S f(z) Q_t^\mu(y, dz) Q_s^\mu(x, dy) \\ &= \int_S \int_S \int_S f(u+v) Q_t(y, du) \mu_t(dv) Q_s^\mu(x, dy) \\ &= \int_S \int_S \int_S \int_S f(u+v) Q_t(w_1 + w_2, du) \mu_t(dv) Q_s(x, dw_1) \mu_s(dw_2) \\ &= \int_S \int_S \int_S \int_S \int_S f(u_1 + u_2 + v) \mu_t(dv) Q_t(w_1, du_1) Q_t(w_2, du_2) \\ &\quad \times Q_s(x, dw_1) \mu_s(dw_2) \\ &= \int_S \int_S \int_S f(u_1 + v) \mu_{s+t}(dv) Q_t(w_1, du_1) Q_s(x, dw_1) \\ &= \int_S \int_S f(u_1 + v) \mu_{s+t}(dv) Q_{s+t}(x, du_1) \\ &= \int_S f(z) Q_{s+t}^\mu(x, dz), \end{aligned}$$

where we have used (5.17), (5.18) and Fubini's theorem. \square

The next result was pointed out to the author by Zenghu Li. It gives a partial converse to Theorem 5.1. Suppose that E has a neutral element e , and

that $Q_t(e, \cdot) = \delta_e(\cdot)$ for all $t \geq 0$. Assume that $(Q_t, t \geq 0)$ is a Markov kernel that has the branching property (5.17), and that $(Q_t^\mu, t \geq 0)$ is defined as in (5.19), with respect to a family of probability measures $(\mu_t, t \geq 0)$.

Theorem 5.2. *If the kernel $(Q_t^\mu, t \geq 0)$ satisfies the Markov property, then $(\mu_t, t \geq 0)$ is a generalised skew-convolution semigroup, as in (5.18).*

Proof. Fix $f \in B_b(E)$. Using the Chapman-Kolmogorov equations, (5.19), (5.17), and then the Chapman-Kolmogorov equations again, we find that

$$\begin{aligned}
& \int_E f(x) \mu_{s+t}(dx) \\
&= \int_E f(x) Q_{s+t}^\mu(e, dx) \\
&= \int_E \int_E f(x) Q_s^\mu(y, dx) Q_t^\mu(e, dy) \\
&= \int_E \int_E \int_E f(u_1 + u_2) Q_s(y, du_1) \mu_s(du_2) Q_t^\mu(e, dy) \\
&= \int_E \int_E \int_E \int_E f(u_1 + u_2) Q_s(y_1 + y_2, du_1) \mu_s(du_2) Q_t(e, dy_1) \mu_t(dy_2) \\
&= \int_E \int_E \int_E \int_E \int_E f(w_1 + w_2 + u_2) Q_s(y_1, dw_1) Q_s(y_2, dw_2) \mu_s(du_2) \\
&\times Q_t(e, dy_1) \mu_t(dy_2) \\
&= \int_E \int_E \int_E \int_E f(w_1 + w_2 + u_2) Q_s(y_1, dw_1) Q_s \mu_t(dw_2) \mu_s(du_2) Q_t(e, dy_1) \\
&= \int_E \int_E \int_E f(w_1 + w_2 + u_2) Q_{s+t}(e, dw_1) Q_s \mu_t(dw_2) \mu_s(du_2) \\
&= \int_E \int_E f(w_1 + x) Q_{s+t}(e, dw_1) (Q_s \mu_t * \mu_s)(dx) \\
&= \int_E f(x) (Q_s \mu_t * \mu_s)(dx),
\end{aligned}$$

and the result follows. \square

There are two other important examples, apart from the Mehler semigroup case:

- Take $E = \mathcal{M}(\mathcal{T})$, the space of all finite Borel measures on a Lusin topological space \mathcal{T} equipped with the topology of weak convergence. The rule $+$ is the usual addition of measures. Then $(Q_t^\mu, t \geq 0)$ is a continuous-state branching process with immigration, where the laws of the immigration process are $(\mu_t, t \geq 0)$. For a monograph account of these processes, see [34]. For interesting connections between catalytic branching processes and Mehler semigroups, see [19]; indeed, in that paper it is shown that Lévy-driven Ornstein-Uhlenbeck processes on $H = L^2((0, \infty))$ arise as limits of fluctuations of the immigration process associated with catalytic branching super-absorbing Brownian motion in $(0, \infty)$.

- Take $E = \mathbb{R}^+ \times \mathbb{R}$, and let $(Q_t, t \geq 0)$ be a homogeneous affine process. In this case $(Q_t^\mu, t \geq 0)$ is a general affine process. This is shown in [18]; applications to finance of general affine processes may be found in [24].

Another, more direct generalisation of (5.13) and (5.14) is to consider two-parameter objects, so instead of a semigroup, we have a two-parameter evolution family of operators $(U(s, t); s \leq t)$ acting in H such that for all $r \leq s \leq t$,

$$U(r, t) = U(s, t)U(r, s),$$

and we seek a family of probability measures $(\mu_{s,t}, s \leq t)$ such that

$$\mu_{r,t} = \mu_{s,t} * (U(s, t)\mu_{r,s}).$$

For a detailed investigation of this set-up, see [38].

6. Invariant measures and operator self-decomposability

Let $Y = (Y(t), t \geq 0)$ be the OU process defined by (3.4), and $(P_t, t \geq 0)$ be the associated Mehler semigroup. We say that $\mu \in \mathcal{M}_1(H)$ is an *invariant measure* for Y if for all $t \geq 0, f \in B_b(H)$,

$$\int_H P_t f(x) \mu(dx) = \int_H f(x) \mu(dx). \tag{6.20}$$

The set of all invariant measures for Y (which may be empty) is convex, and the extremal points are the ergodic measures for Y (see Chapter 3 of [21]). Hence if Y has a unique invariant measure, then it is ergodic. It is well-known that if an invariant measure μ exists for Y , and if we choose Y_0 to have law μ , then Y is a (strictly)-stationary Markov process. In which case, for all $t \geq 0$,

$$Y_0 \stackrel{d}{=} Y(t) = S(t)Y_0 + \int_0^t S(t-r)dL(r),$$

or equivalently,

$$\mu = S(t)\mu * \mu_t, \tag{6.21}$$

where we recall that μ_t is the law of $\int_0^t S(r)dL(r)$.

Measures that have the property (6.21) are called *operator self-decomposable*, and simply *self-decomposable* in the case where $S(t) = e^{-\lambda t}$ for some $\lambda > 0$; we will have more to say about these later on in this section. For now we quote the following from [4], where it appears as Theorem 2.1:

Theorem 6.1. *The following are equivalent:*

1. μ is an invariant measure for Y ;
2. The OU process Y is strictly stationary with $\mathcal{L}(Y_0) = \mu$;
3. μ is operator self-decomposable in the sense of (6.21).

Heuristically, when they exist, invariant measures for a Markov process typically arise as weak limits of the law of the process, corresponding to the dynamical system “settling down after a suitably large time has passed”. We will make

this intuition more precise below; however observe that (6.21) is the formal limit of (4.11), as $t \rightarrow \infty$.

6.1. Invariant measures

A very detailed analysis of invariant measures for Y was carried out by Chojnowska-Michalik [15] (see also [25, 3]). We state her most general results in two theorems:

Theorem 6.2. *If $(\mu_t, t \geq 0)$ converges weakly as $t \rightarrow \infty$, then the limit μ_∞ is an invariant measure for Y . Furthermore, any other invariant measure for Y is of the form $\beta * \mu_\infty$, where $\beta \in \mathcal{M}_1(H)$ is such that $\beta = S(t)\beta$, for all $t \geq 0$.*

Theorem 6.3. *The net $(\mu_t, t \geq 0)$ converges weakly if and only if the following three conditions are satisfied:*

1. $\int_0^\infty \text{tr}(S(t)QS(t)^*)dt < \infty$;
2. $\int_0^\infty \int_H (\|S(t)\|^2 \wedge 1)\nu(dx)dt < \infty$;
3. $\lim_{t \rightarrow \infty} \int_0^t (S(r)b + \int_H S(r)x[\mathbf{1}_{B_1}(S(r)x) - \mathbf{1}_{B_1}(x)]\nu(dx))dr$ exists in H .

The measure μ_∞ , defined in Theorem 6.2, is infinitely divisible, and its characteristics are $(b_\infty, Q_\infty, \nu_\infty)$ where

$$b_\infty = \lim_{t \rightarrow \infty} \int_0^t \left(S(r)b + \int_H S(r)x[\mathbf{1}_{B_1}(S(r)x) - \mathbf{1}_{B_1}(x)]\nu(dx) \right) dr,$$

$$Q_\infty = \int_0^\infty S(t)QS(t)^*dt,$$

$$\nu_\infty(B) = \int_0^\infty \nu_0(S_r^{-1}(B))dr$$

for each $B \in \mathcal{B}(H)$.

There are two special cases of interest that are also treated in [15]:

- If the semigroup $(S(t), t \geq 0)$ is *stable*, i.e. $\lim_{t \rightarrow \infty} S(t)x = 0$, for all $x \in H$, then an invariant measure μ_∞ exists (and is unique) if and only if $\lim_{t \rightarrow \infty} \int_0^t S(r)dL(r)$ converges in probability, in which case μ_∞ is the law of this limiting random variable.
- If the semigroup $(S(t), t \geq 0)$ is *exponentially stable*, i.e. $\|S(t)\| \leq Ce^{-at}$ for some $C \geq 1, a > 0$, then a sufficient condition for existence of a (unique) invariant measure is that

$$\int_H \log^+(\|x\|)\nu(dx) < \infty.$$

In the case where $(S(t), t \geq 0)$ is a contraction semigroup having a bounded generator for which $\lim_{t \rightarrow \infty} \|S(t)\| = 0$, then this logarithmic integrability condition is both necessary and sufficient, as is shown in Theorem 2.1 of [28]. It is well-known that the condition is both necessary and sufficient in finite dimensions (see [5] p. 242, and references therein).

6.2. Operator self-decomposability

Let $\mu \in \mathcal{M}_1(H)$ be arbitrary, and define the set:

$$\mathbb{D}(\mu) = \{T \in L(H); \mu = T\mu * \mu_T \text{ for some } \mu_T \in \mathcal{M}_1(H)\}.$$

The set $\mathbb{D}(\mu)$ is a (algebraic) semigroup with respect to composition of operators; indeed, if $S, T \in \mathbb{D}(\mu)$, then $ST \in \mathbb{D}(\mu)$ with

$$\mu_{ST} = \mu_S * S\mu_T. \quad (6.22)$$

It is shown in [29] that $\mathbb{D}(\mu)$ is closed in the strong topology on $L(H)$. It is called the *Urbanik semigroup* of the measure μ . It is easy to see that μ is operator self-decomposable in the sense of (6.21), if and only if $\mathbb{D}(\mu)$ contains a C_0 -semigroup $(S(t), t \geq 0)$, in which case the measure $\mu_t = \mu_{S(t)}$. Then comparing with (4.11), we see that (6.22) is just the skew-convolution property, hence we may associate a Mehler semigroup with μ , and also an OU process by the construction of section 5, if the semigroup is sufficiently well-behaved.

Operator self-decomposable measures arose in the investigation by Urbanik [53] of certain normalised sums of independent random variables that are associated with uniformly infinitesimal triangular arrays of probability measures on H . The one-dimensional problem was solved by Paul Lévy (see e.g. Theorem 15.3 in [50], p. 51), and gives rise to the classical notion of self-decomposable distribution. Urbanik showed that, under certain technical conditions, operator self-decomposable distributions arise as such limits if and only if $\lim_{t \rightarrow \infty} \|S(t)\| = 0$. Under these conditions he was also able to prove infinite divisibility of the limiting laws. It remains an open problem to extend his results to general C_0 -semigroups. More recently, it has been shown [12] that self-decomposable distributions arise as limits of triangular arrays of strongly mixing sequences of random variables. For more about operator self-decomposability, see [29], and references therein.

If we have an invariant measure μ for the OU semigroup, it is natural to treat μ as a “reference measure”, and to investigate the analytic properties of P_t in the Banach spaces $L^p(H, \mu)$, where $1 \leq p \leq \infty$. Indeed it is easy to see that $(P_t, t \geq 0)$ acts as a contraction semigroup in each of these spaces. It is also a legitimate C_0 -semigroup when $1 \leq p < \infty$ (see the discussion on p. 300 of [32]). A useful tool is the *second quantised* representation of P_t , given by its action on the chaotic decomposition of $L^2(H, \mu)$ into spaces of multiple Wiener-Lévy integrals. In the non-Gaussian case, this was obtained in [41], and from a more general viewpoint in [7, 8]. In [39], the authors extend the work of [49] to establish a Harnack inequality for the Mehler semigroup associated to the OU process. This then implies the strong Feller property for the semigroup, from which we can deduce that the transition probabilities of the process are all absolutely continuous with respect to μ .

7. Cylindrical Ornstein-Uhlenbeck processes

Infinite-dimensional processes arise naturally in mathematical modelling through noise that is described, for each $t \geq 0$ by the formal series:

$$L(t) = \sum_{n=1}^{\infty} \beta_n L_n(t) e_n, \quad (7.23)$$

where $\beta_n \in \mathbb{R}$ for each $n \in \mathbb{N}$, $(L_n, n \in \mathbb{N})$ is a sequence of independent real-valued Lévy processes, and $(e_n, n \in \mathbb{N})$ is a complete orthonormal basis for H . In fact from now on, we will identify H with the space $l^2(\mathbb{N})$ of all real-valued, square-summable sequences. Such a process L is an example of a *cylindrical Lévy process*, and we give a general definition of this concept below. To see that L is not, in general, an H -valued Lévy process, take $\beta_n = 1$ and L_n to be a standard Brownian motion for all $n \in \mathbb{N}$. Then we can easily compute that X has characteristics $(0, I, 0)$; but since the identity I is not trace-class, we do not have a legitimate Brownian motion, or even a Lévy process.

A systematic theory of cylindrical Lévy processes was developed in [6]. This was founded on the theory of cylindrical measures which was developed in the 1960s and 70s by Laurent Schwartz, among others (see e.g. [52]). For $a_1, \dots, a_n \in H$, define $\pi_{a_1, \dots, a_n} : H \rightarrow \mathbb{R}^n$, by $\pi_{a_1, \dots, a_n}(x) = (\langle x, a_1 \rangle, \dots, \langle x, a_n \rangle)$. Then a *cylindrical probability measure* μ is a set function defined on the algebra generated by all cylinder sets of the form $\pi_{a_1, \dots, a_n}^{-1}(B)$ where $B \in \mathcal{B}(\mathbb{R}^n)$, $a_1, \dots, a_n \in H$ and $n \in \mathbb{N}$ such that,

1. $\mu(H) = 1$,
2. The restriction of μ to the σ -algebra generated by $\{\pi_{a_1, \dots, a_n}^{-1}(B), B \in \mathcal{B}(\mathbb{R}^n)\}$ is a bona fide probability measure, for each fixed $a_1, \dots, a_n \in H$ and $n \in \mathbb{N}$.

In general, a *cylindrical process* X is defined to be a family of linear mappings, $(X(t), t \geq 0)$, from H to $L^0(\Omega, \mathcal{F}, P)$. Then for each $t \geq 0$ we obtain a cylindrical probability measure μ_t which plays the role of the ‘‘cylindrical law’’ of $X(t)$ as follows:

$$\mu_t(\pi_{a_1, \dots, a_n}^{-1}(B)) = P((X(t)a_1, \dots, X(t)a_n) \in B).$$

We say that a cylindrical process X is a *cylindrical Lévy process* if for all $t \geq 0$, $a_1, \dots, a_n \in H$, $n \in \mathbb{N}$, $(X(t)a_1, \dots, X(t)a_n)$ is a Lévy process in \mathbb{R}^n . As shown in Lemma 4.2 of [47], the representation (7.23) gives a specific class of examples of this more general notion when we identify $L(t)$ therein with the mapping which sends each $a \in H$ to the random variable $\sum_{n=1}^{\infty} \beta_n L_n(t) \langle e_n, a \rangle$, where the β_n 's are chosen to ensure that the series converges for all $t \geq 0$. In particular, we obtain a cylindrical Lévy process in the case where the L_n 's are i.i.d. centred, square-integrable Lévy processes, and the sequence $(\beta_n, n \in \mathbb{N})$ is bounded.

A version of the Lévy-Itô decomposition holds for cylindrical Lévy processes, and this is used to define stochastic integrals. We may then consider *cylindrical*

Ornstein-Uhlenbeck processes

$$Y(t) = S(t)Y_0 + \int_0^t S(t-r)dX(r),$$

as solutions of the SDE,

$$dY(t) = AY(t)dt + dX(t), \quad (7.24)$$

where the meaning of the last display is precisely that for all $u \in \text{Dom}(A^*)$, $t \geq 0$, with probability 1,

$$Y(t)(u) = Y_0(u) + \int_0^t Y(r)(A^*u)dr + X(t)(u).$$

This is, of course, a natural generalisation to the cylindrical context of the notion of weak solution, as defined in (3.8). Note that $(Y(t), t \geq 0)$ is, in general, itself a cylindrical process; indeed it is cylindrical Markov in that $(Y(t)a_1, \dots, Y(t)a_n, t \geq 0)$ is a Markov process in \mathbb{R}^n for each $a_1, \dots, a_n \in H$, $n \in \mathbb{N}$. Cylindrical notions of Mehler semigroup, invariant measure and self-decomposable distribution are all developed in [6].

When working with cylindrical measures and processes, the notion of *radonification* is important. This refers to finding a mapping into a possibly larger space that transforms the cylindrical object into a bona fide one. An important theorem [26] states that if $(M(t), t \geq 0)$ is a cylindrical semimartingale in H , then there exists a Hilbert-Schmidt operator T on H and a semimartingale $(N(t), t \geq 0)$ so that the real-valued processes $(M(t)(T^*x), t \geq 0)$ and $(\langle N(t), x \rangle, t \geq 0)$ are indistinguishable, for all $x \in H$. In Theorem 5.10 of [47], conditions are found for a suitable deterministic function f so that its cylindrical stochastic integral $\int_0^t f(s)dL(s)$ has the property of *stochastic integrability* in that there exists a random variable I_t such that for all $u \in H$, $\langle I_t, u \rangle = (\int_0^t f(s)dL(s))(u)$ (see also Corollary 4.4 in [46]).

An alternative approach has been developed in a series of papers that focus on the specific class of cylindrical Lévy processes defined by (7.23), with the assumption that the L_n 's are i.i.d. For example [14] considers a Lévy noise obtained by subordinating a cylindrical Wiener process, [44] deals with the case where L_n is a symmetric stable process, while [43] takes L_n to be a symmetric pure jump Lévy process. In the account we give now, we follow [43], but make a slight extension so that L_n is a general symmetric Lévy process (i.e. we include a Gaussian component). So for all $y \in \mathbb{R}$, $n \in \mathbb{N}$, $t \geq 0$,

$$\mathbb{E}(e^{iyL_n(t)}) = \exp \left\{ t \left(-\frac{1}{2}\sigma^2 y^2 + \int_{\mathbb{R}} (\cos(yx) - 1)\nu(dx) \right) \right\}, \quad (7.25)$$

where $\sigma \geq 0$ and ν is a symmetric Lévy measure. Then $L(t)$ takes values in H for all $t \geq 0$ (with probability one) if and only if

$$\sum_{n=1}^{\infty} \left\{ \beta_n^2 \left(1 + \int_{|x| < 1/\beta_n} x^2 \nu(dx) \right) + \int_{|x| \geq 1/\beta_n} \nu(dx) \right\} < \infty.$$

In order to make sense of Ornstein-Uhlenbeck processes within this context, it is necessary to assume that the operator $S(t)$ is compact and self-adjoint for all $t > 0$, and that $(e_n, n \in \mathbb{N})$ in (7.23) is the complete orthonormal basis of eigenvectors that is guaranteed by the Hilbert-Schmidt theorem. It follows that for each $n \in \mathbb{N}, t \geq 0$

$$S(t)e_n = e^{-t\gamma_n} e_n,$$

where $\gamma_n > 0$ and $\lim_{n \rightarrow \infty} \gamma_n = \infty$. We now consider the infinite-dimensional Langevin equation (7.24) as an infinite sequence of one-dimensional equations

$$dY_n(t) = -\gamma_n Y_n(t)dt + \beta_n dL_n(t),$$

having solution starting at $Y(0) = y = (y_n) \in H$, given by

$$Y_n(t) = e^{-t\gamma_n} y_n + \beta_n \int_0^t e^{-\gamma_n(t-s)} dL_n(s).$$

Then $Y(t) = (Y_n(t), n \in \mathbb{N})$ is a bona-fide stochastic process taking values in H for all $t \geq 0$ (with probability one) if and only if

$$\sum_{n=1}^{\infty} \frac{\beta_n^2}{\gamma_n} + \sum_{n=1}^{\infty} \frac{1}{\gamma_n} \int_{1/\beta_n}^{e^{\gamma_n}/\beta_n} \left(\frac{\psi_0(u)}{u^3} + \frac{\psi_1(u)}{u} \right) du < \infty, \quad (7.26)$$

where $\psi_0(u) := \int_{|x| \leq u} x^2 \nu(dx)$ and $\psi_1(u) = \int_{|x| > u} \nu(dx)$. Furthermore, the process is adapted and Markov. In particular, (7.26) is satisfied if the sequence (β_n) is bounded, and the following holds:

$$\int_1^{\infty} \log(y) \nu(dy) < \infty \text{ and } \sum_{n=1}^{\infty} 1/\gamma_n < \infty,$$

in which case the process has a unique invariant measure. In Corollary 6.3 of [47], it is shown that (7.26) is a necessary and sufficient condition for the semigroup $(S(t), t \geq 0)$ to be stochastically integrable, (without the assumption that the L_n s are identically distributed).

The investigation of càdlàg paths in this context has attracted some attention. In [13], the authors show that if the vectors $(e_n, n \in \mathbb{N})$ are all in $\text{Dom}(A^*)$ and the sequence (β_n) fails to converge to zero, then with probability 1, the OU process Y has no point $t > 0$ at which either the right or left limit exists. In [45], conditions are found on the Lévy measure for Y to have a càdlàg modification, provided the semigroup $(S(t), t \geq 0)$ is both exponentially stable and analytic.

In the case where the L_n 's have symmetric α -stable distributions ($0 < \alpha < 2$), a finer analysis may be carried out; in [44] the authors obtain a number of interesting properties of the solution, including gradient estimates on the transition semigroup, from which they deduce the strong Feller property; while in [35] it is shown that the OU process Y has a càdlàg modification if and only if $\sum_{n=1}^{\infty} |\beta_n|^\alpha < \infty$.

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